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Cite as: Appl. Phys. Lett. **120**, 051901 (2022); https://doi.org/10.1063/5.0074419 Submitted: 08 October 2021 • Accepted: 13 January 2022 • Published Online: 04 February 2022

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## Epitaxial growth and thermoelectric properties of Mg<sub>3</sub>Bi<sub>2</sub> thin films deposited by magnetron sputtering (1)

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Note: This paper is part of the APL Special Collection on Thermoelectric Materials Science and Technology Towards Applications.

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#### **ABSTRACT**

 $Mg_3Sb_2$ -based thermoelectric materials attract attention for applications near room temperature. Here, Mg-Bi films were synthesized using magnetron sputtering at deposition temperatures from room temperature to  $400\,^{\circ}$ C. Single-phase  $Mg_3Bi_2$  thin films were grown on c-plane-oriented sapphire and Si(100) substrates at a low deposition temperature of  $200\,^{\circ}$ C. The  $Mg_3Bi_2$  films grew epitaxially on c-sapphire and fiber-textured on Si(100). The orientation relationships for the  $Mg_3Bi_2$  film with respect to the c-sapphire substrate are (0001)  $Mg_3Bi_2\|(0001)$   $Al_2O_3$  and  $[11\overline{2}0]$   $Mg_3Bi_2\|[11\overline{2}0]$   $Al_2O_3$ . The observed epitaxy is consistent with the relatively high work of separation, calculated by the density functional theory, of 6.92 J m $^{-2}$  for the  $Mg_3Bi_2$  (0001)/ $Al_2O_3$  (0001) interface.  $Mg_3Bi_2$  films exhibited an in-plane electrical resistivity of  $34\,\mu\Omega$  m and a Seebeck coefficient of  $+82.5\,\mu$ V K $^{-1}$ , yielding a thermoelectric power factor of  $200\,\mu$ W m $^{-1}$  K $^{-2}$  near room temperature.

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Thermoelectric thin films are desired for microscale devices using semiconductor coolers<sup>1</sup> or self-powered systems for sensors and the Internet of Things.<sup>2,3</sup> The interest in thermoelectric thin films is also largely driven by possible improvement in thermoelectric properties, for example, in lower dimensional structures due to quantum confinement<sup>4</sup> and phonon-scattering effects.<sup>5</sup> Compared with bulk materials, thermoelectric thin films, especially single-crystal thin films, are attractive for assembling micro-thermoelectric generators.<sup>2,3</sup>

Conventional thermoelectric materials, most notably  $Bi_2Te_3$ -based ones, have some well-known disadvantages in that they are brittle and contain the extremely scarce element Te. The Zintl phase  $Mg_3Bi_2$  and its alloys in a  $La_2O_3$ -type trigonal structure have received attention because of their nontoxicity, high elemental abundancy, and good thermoelectric properties for energy harvesting applications. <sup>6–15</sup>

For example, an average thermoelectric figure of merit of polycrystal-line  $\mathrm{Mg_{3+\delta}Sb_xBi_{2-x}}$  has been reported to be comparable with that of the conventional n-type  $\mathrm{Bi_2Te_{2.7}Se_{0.3}}$  in the temperature range of 50–250 °C. <sup>10</sup> Further studies have focused on the improvement of thermoelectric properties near room temperature. <sup>1,11,12,14,16,17</sup> Single-crystal  $\mathrm{Mg_3Bi_2}$  and its alloys were synthesized by the Mg-flux method. <sup>8,13</sup> For vacuum-based thin-film deposition, the relatively high vapor pressure of Mg (of the order of  $10^{-6}$  mbar at 250 °C and  $10^{-4}$  mbar at 330 °C) <sup>18</sup> makes it challenging to synthesize pure  $\mathrm{Mg_3Bi_2}$  thin films at high temperatures. Matsui et al. reported growing polycrystal-line  $\mathrm{Mg_3Bi_2}$  thin films on a copper plate at 200 °C by radio frequency magnetron sputtering. <sup>19</sup> Single-crystalline  $\mathrm{Mg_3Bi_2}$  thin films deposited using molecular beam epitaxy have been reported. <sup>20,21</sup> However, none of the above studies were aiming for thermoelectric applications.

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The growth of epitaxial thin films is, therefore, important to enable characterization and explanation of the fundamental physical properties of  $\rm Mg_3Bi_2$ -based thin films. Here, we grow a  $\rm Mg_3Bi_2$  thin film on c- and r-plane-oriented sapphire and  $\rm Si(100)$  substrates by magnetron co-sputtering of Mg and Bi targets. Epitaxial single-phase  $\rm Mg_3Bi_2$  is synthesized on the c-sapphire substrate, and a fiber-textured  $\rm Mg_3Bi_2$  film on  $\rm Si(100)$  substrates. The orientation relationships for  $\rm Mg_3Bi_2$  with respect to the c-sapphire substrate are: (0001)  $\rm Mg_3Bi_2$  with respect to the c-sapphire substrate are: (0001)  $\rm Mg_3Bi_2$  ||(0001)  $\rm Al_2O_3$  and [11 $\overline{2}$ 0]  $\rm Mg_3Bi_2$ ||[11 $\overline{2}$ 0]  $\rm Al_2O_3$ . Moreover, the in-plane thermoelectric properties of the  $\rm Mg_3Bi_2$  films grown on c-sapphire were measured from room temperature to deposition temperatures.

Mg<sub>3</sub>Bi<sub>2</sub> films were deposited using direct current (dc) magnetron sputtering in an ultrahigh vacuum chamber (base pressure  $< 4 \times 10^{-6}$ Pa). The deposition system is described elsewhere. 22 Mg and Bi circular targets (50 mm in diameter) were used and driven by dc-power supplies with optimized (i.e., yielding approximately the desired 3:2 composition; see the supplementary material, Fig. S1) power of 90 and 15 W, respectively. The Ar gas flow rate was fixed to 80 sccm corresponding to a working pressure of 0.5 Pa. The substrate holder was maintained at different temperatures from non-intentional heating ("room temperature," RT) to 400 °C, rotated at 15 rpm, with the substrates electrically floating. The films were deposited on Si(100) and c- and r-plane-oriented sapphire substrates. The films deposited on Si(100) and c-sapphire substrates were used for microstructural characterization by x-ray diffraction (XRD) and elemental composition determined by energy-dispersive x-ray spectrometry (EDS). The electrical resistivity  $\rho$ , the Seebeck coefficient S, and the power factor  $S^2/\rho$  of films on c- and r-sapphire were measured by a commercial system (CTA-3, Cryoall Co., Ltd., China) under a low-pressure helium atmosphere from RT to 200 °C (growth temperature). The uncertainty for thermoelectric property measurements is around 5%. The Hall coefficient R<sub>H</sub> for Mg<sub>3</sub>Bi<sub>2</sub> grown on c-sapphire was obtained from the magnetic field-dependent Hall resistivities using the electrical transport option in a quantum design physical property measurement

system (PPMS). The charge carrier concentration  $n_{\rm H}$  was calculated by  $n_{\rm H} = 1/eR_{\rm H}$ . The work of separation<sup>23</sup> was calculated at 0 K for the interface of interest using the density functional theory (DFT) as implemented in the Vienna *ab initio* simulation package (VASP). <sup>24–26</sup> The generalized gradient approximation was employed in VASP and parametrized by Perdew, Burke, and Ernzerhof. <sup>27</sup> The precision and convergence for the interface were specified by a reciprocal k-grid of 5  $\times$  5  $\times$  1, 520 eV cutoff, and 0.01 meV criterion for the total energy. Full structural optimization was made by minimizing the interatomic forces and optimizing the lattice parameters.

Figure 1 shows XRD patterns of as-deposited Mg-Bi films on c-sapphire and Si(100) with  $T_s$  varied from RT to 400 °C. The elevated growth temperature improved the crystallinity of Mg-Bi films on both substrates. For the films deposited on c-sapphire, the peaks of the Mg-Bi film deposited at RT can be attributed to the polycrystalline Mg<sub>3</sub>Bi<sub>2</sub> phase. When increasing T<sub>s</sub> to 200 °C, a 0001 peak of Mg<sub>3</sub>Bi<sub>2</sub> shows up at 11.9°, and the intensities of corresponding higher-order 000l peaks increase, indicating the possibility of epitaxial growth. Films deposited at higher  $T_s$  have clearer, sharper peaks corresponding to Mg<sub>3</sub>Bi<sub>2</sub>, but also exhibit an additional peak located at  $27.0^{\circ}$ , which can be attributed to the  $01\overline{1}2$  peak of Bi. The presence of metallic Bi is correlated with Mg deficiency for deposition at higher temperatures ( $T_s \ge 300\,^{\circ}\text{C}$ , see the supplementary material, Fig. S1). The films grown on Si(100) at  $T_s \le 200$  °C show the same Mg<sub>3</sub>Bi<sub>2</sub> phase with similar orientations with films grown on c-sapphire, while the films grown with  $T_s > 200$  °C exhibit  $11\overline{2}2$ ,  $11\overline{2}3$ ,  $10\overline{1}5$ , and  $12\overline{3}4$  peaks of substantial intensity, indicating the presence of other orientations. Therefore,  $T_s = 200 \,^{\circ}\text{C}$  is an optimal temperature, sufficient to grow single-phase Mg<sub>3</sub>Bi<sub>2</sub> on both substrates, which is in agreement with the previous work<sup>28</sup> on the synthesis of Mg<sub>3</sub>Bi<sub>2</sub> films using radio frequency magnetron sputtering.

To further identify the crystallographic relationship between Mg<sub>3</sub>Bi<sub>2</sub> films and substrates, pole figures of (0001) and (11 $\overline{2}$ 2) planes were acquired for Mg<sub>3</sub>Bi<sub>2</sub> films grown at  $T_s = 200\,^{\circ}\text{C}$  on c-sapphire and Si(100) substrates, as shown in Fig. 2. The pole figure of the 0001

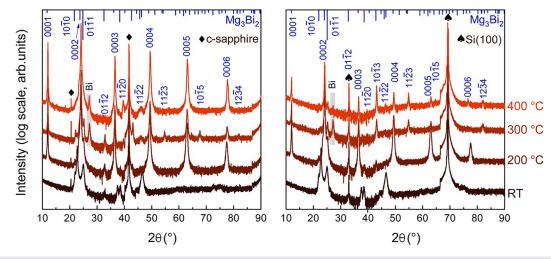


FIG. 1. X-ray diffraction patterns of Mg<sub>3</sub>Bi<sub>2</sub> thin films grown on c-sapphire and Si(100) at different temperatures. The peak position of the Mg<sub>3</sub>Bi<sub>2</sub> phase is from powder diffraction file No. 65-8732.

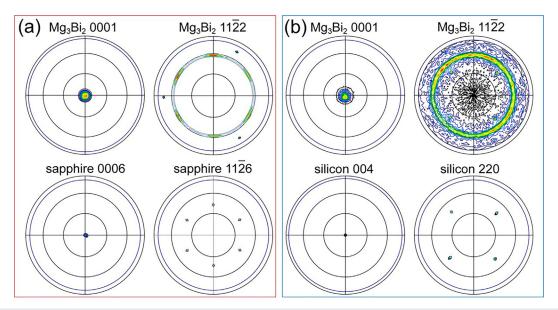


FIG. 2. 0001 and  $11\overline{2}2$  pole figures of Mg<sub>3</sub>Bi<sub>2</sub> thin films deposited at 200 °C on the c-sapphire substrate (a) and Si(100) substrate (b), and corresponding pole figure 0006 and  $11\overline{2}6$  of c-sapphire, 004 and 220 of the silicon substrate.

peak [Fig. 2(a), top left] shows diffraction only at the center (i.e., tiltangle  $\chi = 0$ ) consistent with basal-plane orientation [Fig. 2(a), bottom left]. The  $11\overline{2}2$  pole figure [Fig. 2(a), top left] shows a set of sixfold symmetric spots out of a visible ring at  $\chi = 57.5^{\circ}$ , corresponding to the angle between (000*l*) and (1122) planes in the trigonal Mg<sub>3</sub>Bi<sub>2</sub> structure. Their in-plane rotation angles  $\varphi$  are matched with the same diffraction peaks of the  $11\overline{2}6$  peak of Al<sub>2</sub>O<sub>3</sub>, consistent with the  $\varphi$  scans (the supplementary material, Fig. S2), indicating that  $Mg_3Bi_2$  (11 $\overline{2}$ 2) planes have the same in-plane orientation as  $Al_2O_3$  (11 $\overline{2}6$ ). We also find a set of threefold symmetric diffraction peaks at  $\chi = 72.4^{\circ}$ , originating from the  $(20\overline{2}2)$  plane of sapphire. Based on these results, we conclude that the orientation relationships for the Mg<sub>3</sub>Bi<sub>2</sub> film with respect to the c-sapphire substrate are: (0001) Mg<sub>3</sub>Bi<sub>2</sub> $\parallel$ (0001) Al<sub>2</sub>O<sub>3</sub> and [11 $\overline{2}$ 0]  $Mg_3Bi_2||[11\overline{2}0] Al_2O_3$ . In comparison, the 0001 and  $11\overline{2}2$  pole figures of Mg<sub>3</sub>Bi<sub>2</sub> films grown on Si(100) are shown in Fig. 2(b), top. The 0001 pole figure shows diffraction only at the center, whereas the 1122 pole figure exhibits a broad ring of uniform intensity, indicating that the majority of grains grow in the same direction out of the plane but are randomly distributed in all directions in-plane, i.e., fiber texture.

To rationalize the epitaxy of  $Mg_3Bi_2$  on sapphire, the interface structure is analyzed. A graphical in-plane view of the interface is shown in Fig. 3(a). The symmetries of the two lattices are close to each other, as confirmed by the in-plane mismatch value of 1.7% between  $Mg_3Bi_2$  and  $Al_2O_3$ .<sup>21</sup> To quantify the interfacial strength, the work of separation was calculated by DFT, as illustrated in Fig. 3(b). A mixture of covalent and ionic bonding between Mg and O as well as Bi and O across the interface, each being threefold coordinated, is observed, implying strong bonding. The work of separation increases, while the interface energy decreases. The work of separation of 6.92 J m<sup>-2</sup> was obtained for the  $Mg_3Bi_2$  (0001)/ $Al_2O_3$  (0001) interface, which is sufficiently strong to facilitate epitaxial growth in comparison with Cu (111)/ $Al_2O_3$  (0001) and  $V_2AlC$  (0001)/ $Al_2O_3$  (11 $\overline{2}$ 0),  $\overline{3}^1$  having the work of separation of 5.48 and

 $2.86\,\mathrm{J}\;\mathrm{m}^{-2}$ , respectively. However, this work of separation is also weaker than  $12.70\,\mathrm{J}\;\mathrm{m}^{-2}$  for Nb  $(111)/\mathrm{Al_2O_3}$   $(0001)^{32}$  and  $9.31\,\mathrm{J}\;\mathrm{m}^{-2}$  for the TiO<sub>2</sub>  $(100)/\mathrm{Al_2O_3}$  (0001) interface. It may lead to the epitaxial quality of Mg<sub>3</sub>Bi<sub>2</sub> on sapphire not being as good as expected with some random reflections observed from Mg<sub>3</sub>Bi<sub>2</sub> in the XRD pattern [Fig. 1(a)] and a visible ring in pole figures [Fig. 2(a)].

Figure 4 shows the temperature-dependent thermoelectric properties of the single-phase Mg<sub>3</sub>Bi<sub>2</sub> films deposited on c- and r-sapphire (both having the same crystalline feature, Fig. S3). The electrical resistivity  $\rho$  of the Mg<sub>3</sub>Bi<sub>2</sub> film is essentially constant with a variation from 34.5 to 36.8  $\mu\Omega$  m (i.e., within error margins) in the measured temperature range between 30 and 200 °C. These obtained values of the resistivity are consistent with the range expected for the semimetal Mg<sub>3</sub>Bi<sub>2</sub> with a small band overlap energy at room temperature. 11,34 The Mg<sub>3</sub>Bi<sub>2</sub> film exhibits p-type behavior and a linearly decreasing Seebeck coefficient from +82.5 to  $+53.4 \,\mu\text{V K}^{-1}$  when temperature increases from RT to 200  $^{\circ}\text{C}.$  The power factor  $(\text{S}^2\!/\!\rho)$  of the  $\text{Mg}_3\text{Bi}_2$  film is 197 at 30 °C and 79  $\mu$ W m<sup>-1</sup> K<sup>-2</sup> at 200 °C. A comparison of thermoelectric properties at 30 °C between the epitaxial Mg<sub>3</sub>Bi<sub>2</sub> film in this work and reported values for single-crystal and bulk Mg3Bi2 materials is shown in the supplementary material, Table S1. There is a significant difference in the electrical resistivity due to anisotropy and the synthesis methods. 13,35 Generally, the in-plane resistivity is higher than that of out-of-plane. The in-plane resistivity of the Mg<sub>3</sub>Bi<sub>2</sub> thin film in this work is 34–36  $\mu\Omega$  m. Kim et al.<sup>36</sup> found a large anisotropy in Mg<sub>3</sub>Bi<sub>2</sub> single crystals. Xin et al.8 recently showed that the out-of-plane resistivity (along the *c* axis) is only half of the out-of-plane one (along the *b* axis) in the temperature range 2-320 K. There are also differences in the measured resistivity of bulk polycrystalline Mg<sub>3</sub>Bi<sub>2</sub> compounds prepared by mechanical alloying<sup>37</sup> and by solid-state reaction.<sup>38</sup> Therefore, the synthesis methods play a role in determining the impurity content, grain sizes, and composition, significantly affecting

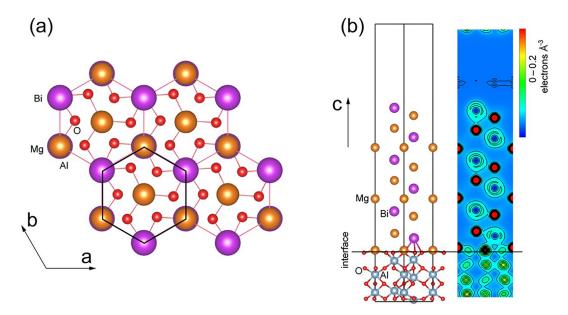
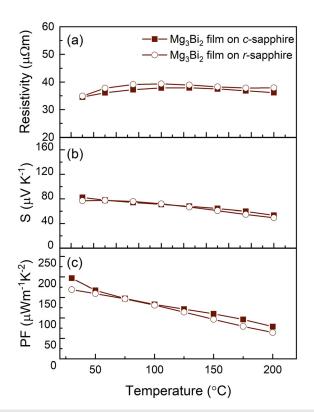


FIG. 3. (a) Ball-stick model of Mg<sub>3</sub>Bi<sub>2</sub> (0001) on O-terminated sapphire as an in-plane view. (b) Structural model of the Mg<sub>3</sub>Bi<sub>2</sub>/Al<sub>2</sub>O<sub>3</sub> interface together with a corresponding electron density distribution.



**FIG. 4.** Thermoelectric properties of the Mg<sub>3</sub>Bi<sub>2</sub> thin film deposited at 200 °C on *c*-sapphire (filled box) and *r*-sapphire (empty circle): electrical resistivity  $\rho$ , Seebeck coefficient (*S*), and power factor  $S^2/\rho$  (PF) as a function of temperature.

electrical transport. Also, the Mg3Bi2 materials mentioned above are all p-type semiconductors. However, with Mg overstoichiometry and Te doping, this phase may turn to *n*-type. The crystalline quality and the presence of dopants can also affect the carrier concentration  $(n_{\rm H})$ and mobility ( $\mu_{\rm H}$ ). In our case (Fig. S4),  $n_{\rm H}$  of the Mg<sub>3</sub>Bi<sub>2</sub> film grown on *c*-sapphire was  $\sim 1.7 \times 10^{19}\,{\rm cm}^{-3}$  at 2 K and slightly increased to  $4.0 \times 10^{19} \, \mathrm{cm}^{-3}$  at 300 K. These values are in good agreement with earlier reports<sup>8,11,39</sup> (several tens of  $10^{19} \, \mathrm{cm}^{-3}$ ). The  $\mu_{\mathrm{H}}$  decreased from 59.1 to  $14.0 \,\mathrm{cm}^2 \,\mathrm{V}^{-1} \,\mathrm{s}^{-1}$  when the temperature increased from 2 to 300 K. This trend is generally due to more frequent carrierscattering by phonons at higher temperatures. Compared to previous results on the Mg<sub>3.2</sub>Bi<sub>2</sub> single-crystal bulk<sup>8</sup> and the Mg<sub>3</sub>Bi<sub>2</sub> film grown by molecular beam epitaxy,  $^{99}$   $\mu_{\rm H}$  of the Mg<sub>3</sub>Bi<sub>2</sub> film in this work is dropping significantly earlier, above 10 K, which may be due to predominant grain-boundary scattering. The behavior of  $\mu_{\rm H}$  of  $n_{\rm H}$  indicates a semimetallic characteristic and can be influenced by crystalline quality, in turn affecting defect scattering of charge carriers. As temperature increases above room temperature, n<sub>H</sub> increases sharply (which is favorable for high S and low  $\rho$ ) while  $\mu_{\rm H}$  decreases (favorable for low S and high  $\rho$ ); the resistivity increases (which is usual for semimetals) while the Seebeck coefficient decreases. The sharp increase in the carrier concentration with temperature could be a sign of the bipolar effect, 40 which is common in semimetals and narrow-gap semiconductors. The bipolar effect is detrimental to the Seebeck effect. 41 A decreasing bipolar effect can be achieved through increasing the majority carrier concentration, for example, through doping, and decreasing the minority carrier concentration.

In conclusion, we have demonstrated that single-phase  $Mg_3Bi_2$  thin films can be grown on different substrates at a low deposition temperature of 200 °C. The  $Mg_3Bi_2$  films grew epitaxially on c-sapphire and fiber-textured on Si(100) substrates. The orientation

relationships for the Mg<sub>3</sub>Bi<sub>2</sub> film with respect to the *c*-sapphire substrate are (0001) Mg<sub>3</sub>Bi<sub>2</sub>||(0001) Al<sub>2</sub>O<sub>3</sub> and [11 $\overline{2}$ 0] Mg<sub>3</sub>Bi<sub>2</sub>||[11 $\overline{2}$ 0] Al<sub>2</sub>O<sub>3</sub>. The observed epitaxy is consistent with the relatively high work of separation, calculated by DFT, of 6.92 J m<sup>-2</sup> obtained for the Mg<sub>3</sub>Bi<sub>2</sub> (0001)/Al<sub>2</sub>O<sub>3</sub> (0001) interface. The in-plane thermoelectric properties of the Mg<sub>3</sub>Bi<sub>2</sub> films grown on *c*-sapphire were measured from room temperature to 200 °C. The in-plane resistivity of Mg<sub>3</sub>Bi<sub>2</sub> films is 34  $\mu$ Ωm near room temperature, and the Seebeck coefficient linearly decreases from 82.5 to 53.4  $\mu$ V K<sup>-1</sup> from room temperature to 200 °C. The power factor is about 200  $\mu$ Wm<sup>-1</sup> K<sup>-2</sup> near room temperature. The work has prospects for preparing Mg<sub>3</sub>Bi<sub>2</sub>-based thin-film thermoelectrics applied near room temperature and future thermoelectric property improvement for Mg<sub>3</sub>Bi<sub>2</sub> through doping.

See the supplementary material for additional XRD and composition data, comparison of thermoelectric properties, Hall carrier concentration, and carrier mobility.

This work was supported financially by the Swedish Government Strategic Research Area in Materials Science on Functional Materials at Linköping University (Faculty Grant SFO-Mat-LiU No. 2009 00971), the Knut and Alice Wallenberg Foundation through the Wallenberg Academy Fellows program (No. KAW-2020.0196), the Swedish Research Council (VR) under Project Nos. 2016-03365 and 2021-03826, the National Key Research and Development Program of China under Grant No. 2018YFB0703600, the National Natural Science Foundation of China under Grant No. 51872133, the Guangdong Innovative and Entrepreneurial Research Team Program under Grant No. 2016ZT06G587, and the Tencent Foundation through the XPLORER PRIZE, Guangdong Provincial Key Laboratory Program (No. 2021B1212040001) from the Department of Science and Technology of Guangdong Province. The computations were performed on resources provided by the Swedish National Infrastructure for Computing (SNIC) at National Supercomputer Centre (NSC) partially funded by the Swedish Research Council through Grant Agreement No. 2018-05973.

### AUTHOR DECLARATIONS Conflict of Interest

The authors have no conflicts to disclose.

#### **Author Contributions**

R.S. and P.E. initiated the study. P.E., W.L. and A.I.F. supervised the research. G.S. and R.S. planned and performed the sample synthesis and microstructure characterization with contributions from A.I.F., and P.E., G.S, and Y.Z. performed and analyzed thermoelectric property measurements. T.F. performed and analyzed the Hall measurements. D.M. performed DFT calculations and wrote the corresponding part of the manuscript. G.S., R.S., and P.E. wrote the manuscript with contributions from the co-authors. All co-authors read, edited, and commented on successive version of the manuscript.

#### DATA AVAILABILITY

The data that support the findings of this study are available within the article and its supplementary material or are available from the corresponding authors upon reasonable request.

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