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Anomally strong pinning of the filling factor $\nu = 2$ in epitaxial graphene

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We explore the robust quantization of the Hall resistance in epitaxial graphene grown on Si-terminated SiC. Uniquely to this system, the dominance of quantum over classical capacitance in the charge transfer between the substrate and graphene is such that Landau levels (in particular, the one at exactly zero energy) remain completely filled over an extraordinarily broad range of magnetic fields. One important implication of this pinning of the filling factor is that the system can sustain a very high nondissipative current. This makes epitaxial graphene ideally suited for quantum resistance metrology, and we have achieved a precision of 3 parts in $10^{10}$ in the Hall resistance-quantization measurements.

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The quantum Hall effect (QHE) is one of the key fundamental phenomena in solid-state physics. It was observed in two-dimensional electron systems in semiconductor materials and, since recently, in graphene: both in exfoliated and epitaxial devices. A direct high-accuracy comparison of the conventional QHE in semiconductors with that observed in graphene constitutes a test of the universality of this effect. The affirmative result would strongly support the pending redefinition of the SI units based on the Planck constant $h$ and the electron charge $e$ and provide an international resistance standard based upon quantum physics.

Graphene is believed to offer an excellent platform for QHE physics due to the large energy separation between Landau levels (LL) resulting from the Dirac-type "massless" electrons specific for its band structure. The Hall resistance quantization with an accuracy of 3 parts in $10^9$ has already been established in Hall-bar devices manufactured from epitaxial graphene grown on Si-terminated face of SiC (SiC/G). However, for graphene to be practically employed as an embodiment of a quantum-resistance standard, it needs to satisfy further stringent requirements, in particular with respect to robustness over a range of temperature, magnetic field, and measurement current. A high measurement current, which a device can sustain at a given temperature without dissipation, is particularly important for precision metrology as it defines the maximum attainable signal-to-noise ratio.

The extent of the QHE plateau in conventional 2D electron systems is, usually, set by disorder and temperature. Disorder pins the Fermi energy in the mobility gap of the 2D system, which suppresses dissipative transport at low temperatures over a finite range of magnetic fields around the values corresponding to exactly filled LLs. These values can be calculated from the carrier density $n_e$ determined from the low-field Hall resistivity measurements and coincide with the maximum nondissipative current, the breakdown current. Thus, the breakdown current in conventional two-dimensional semiconductors peaks very close to the field values where the filling factor $\nu$ is an even integer. Though less studied experimentally, the behavior of the breakdown current on the plateau for the exfoliated graphene, including the $\nu = 2$ plateau corresponding to the topologically protected $N = 0$ LL, looks quite similar.

In this Brief Report, we explore the robustness of the Hall resistance quantization in SiC/G. Unlike the QHE in conventional 2D systems, where the carrier density is independent of magnetic field, here specifically to SiC/G, we find that the carrier density in graphene varies with magnetic field due to the charge transfer between surface-donor states in SiC and graphene. Most importantly, we find magnetic field intervals of several Tesla, where the carrier density in graphene increases linearly with the magnetic field, resulting in the pinning of $\nu = 2$ state with electrons at the chemical potential occupying SiC surface donor states half-way between the $N = 0$ and $N = 1$ LLs in graphene. Interestingly, at magnetic fields above the $\nu = 2$ filling-factor pinning interval, the carrier density saturates at a value up to 30% higher than the zero-field carrier density. The pinned filling factor manifests itself in a continuously increasing breakdown current toward the upper magnetic field end of the $\nu = 2$ state far beyond the nominal value of $B_{\nu=2}$ calculated from the zero-field carrier density. Facilitated by the high breakdown current in excess of 500 $\mu$A at 14 T, we have achieved a precision of 3 parts in $10^{10}$ in the Hall resistance quantization measurements.

The anomalous pinning of $\nu = 4N + 2$ filling factors in SiC/G is determined by the dominance of the quantum capacitance, $c_q$, over the classical capacitance per unit area, $c_c$, in the charge transfer between graphene and surface-donor states of SiC/G: $c_q \ll c_c$, where $c_q = e^2 \gamma e$, $c_c = 1/(4\pi d)$, and $\gamma e$ is the density of states of electrons at the Fermi level. The latter reside in the “dead layer” of carbon atoms, just underneath graphene. This layer is characterized by a $6\sqrt{3} \times 6\sqrt{3}$ supercell of the reconstructed surface of sublminated SiC. Missing or substituted carbon atoms in various positions of such a huge supercell in the dead layer create localized surface states with a broad distribution of energies within the bandgap of SiC ($\sim 2.4$ eV).

It appears that the density of such defects is higher in material grown at low temperatures (1200 – 1600°C), resulting in graphene doped to a large electron density, $n_e \sim 10^{13}$ cm$^{-2}$, which is difficult to change. On the other hand,
FIG. 1. (Color online) Schematic band-structure for graphene on SiC in zero field (a) and in quantizing field (b); graphical solution for carrier density as a function of magnetic field, $n_s(B)$, of the charge-transfer model given by Eq. (1) (black line) together with lines of constant filling factor (red/gray lines) and for carrier density as a function of magnetic field, SiC in zero field (a) and in quantizing field (b); graphical solution ng

The quantum capacitance of a two-dimensional electron system is the result of a low compressibility of the electron liquid determined by the peaks in $\gamma_e$. For electrons in high-mobility GaAs/AlGaAs heterostructures in magnetic field, the quantum capacitance manifests itself in weak magneto-oscillations of the electron density due to the suppressed density of states inside the inter-Landau level gaps. A similarly weak effect has been observed in graphene exfoliated onto n-Si/SiO$_2$ substrate, where the influence of a larger (than in usual semiconductors) inter-LL gaps is hindered by a strong charging effect determined by a relatively large thickness of SiO$_2$ layer. For epitaxial graphene on SiC, due to the short charging effect determined by a relatively large thickness of usual semiconductors) inter-LL gaps is hindered by a strong density of states inside the inter-Landau level gaps. A similarly take the form of the robust pinning of the electron-filling system lies inside the gap between $\tilde{\nu}$ and $\tilde{\nu}=\tilde{\nu}_{\text{ll}}$, of the visible only in the case of the higher $n_s$ in graphene, the effect of quantum capacitance manifests itself in weak magneto-oscillations of the electron density due to the suppressed density of states inside the inter-Landau level gaps. A similarly weak effect has been observed in graphene exfoliated onto n-Si/SiO$_2$ substrate, where the influence of a larger (than in usual semiconductors) inter-LL gaps is hindered by a strong charging effect determined by a relatively large thickness of SiO$_2$ layer. For epitaxial graphene on SiC, due to the short charging effect determined by a relatively large thickness of usual semiconductors) inter-LL gaps is hindered by a strong density of states inside the inter-Landau level gaps. A similarly...
Determination of the breakdown current, $I_c$, quantum Hall resistivity values for filling factors $\nu$. The horizontal lines indicate the exact measurement of $\rho_{xy}$ configuration explained in legend. (c) High-precision GaAs quantum Hall sample.7 All error bars are $1/\sqrt{N}$, where $N = 1000$. The as-grown sample had the zero-field carrier density $n_s = 1.1 \times 10^{12} \text{ cm}^{-2}$. Graphene was encapsulated in a polymer bilayer, a spacer polymer followed by an active polymer able to generate acceptor levels under UV light. At room temperature, electrons diffuse from graphene through the spacer polymer layer and fill the acceptor levels in the top polymer layer. Such a photochemical gate allowed nonvolatile control over the charge-carrier density in graphene. More fabrication details can be found elsewhere.7,22

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Figure 2(a) shows magneto transport measurements on the encapsulated sample tuned to a zero-field carrier density of $n_s = 6.7 \times 10^{11} \text{ cm}^{-2}$ corresponding to the case in Fig. 1(c). From the carrier density we estimate that the magnetic field $B_{\nu=2}$ needed for exact filling factor $\nu = 2$ in this device is 13.8 T. A well-quantized Hall plateau in $\rho_{xy}$ can be seen at $\nu = \pm 2$ for both magnetic field directions, which is more than 5 T wide, whereas the longitudinal resistivity, $\rho_{xx}$, drops to zero signifying a nondissipative state. In addition, a less precisely quantized plateau is present at $\nu = \pm 6$, for which $\rho_{xx}$ remains finite.

Accurate quantum Hall resistance measurements require that the longitudinal voltage remains zero (in practice, below the noise level of the nanovolt meter) to ensure the device is in the nondissipative state, which can be violated by the breakdown of the QHE at high current. Figure 2(b) shows the determination of the breakdown current $I_c$ at $B = 14 \text{ T}$ on the $\nu = 2$ plateau. Here we define $I_c$ as the source-drain current, $I_{sd}$, at which $V_{xx} > 10 \text{ nV}$. We find for three different combinations of source-drain current contacts that the breakdown current for this value of $n_s$ is approximately $50 \mu\text{A}$ (note that $I_{sd}$ in a practical quantum Hall to 100 $\Omega$ resistance measurement is $\approx 25 \mu\text{A}^{27}$). The contact resistance, determined via a three-terminal measurement in the nondissipative state, is smaller than 1.5 $\Omega$.

Figure 2(c) shows a precision measurement of $\rho_{xy}$ and $\rho_{xx}$ for different magnetic fields along the $\nu = 2$ plateau. Note that this plateau appears much shorter in the magnetic-field range than that shown in Fig. 2(a) because of the 200 times larger measurement current used in precision measurements. From this figure we determine that the mean of $\Delta \rho_{xy}/\rho_{xy}$ is $0.06 \pm 0.3 \times 10^{-9}$ for the data between 11.75 and 14.0 T, while at the same time $\rho_{xx} < 1 \text{ m}\Omega$. This result represents an order of magnitude improvement of QHE precision measurements in graphene, as compared to the earlier record.7 Not only is QHE accurate, but it is also extremely robust in this epitaxial graphene device, easily meeting the stringent criteria for accurate quantum Hall resistance measurements normally applied to semiconductor systems.

Using the polymer gating method,22 we further reduce the zero-field electron density $n_s$ in graphene to correspond to the solution of the charge transfer problem in Fig. 1(d), i.e., down to $4.6 \times 10^{11} \text{ cm}^{-2}$ as evidenced by magnetotransport measurements in Fig. 3(a). On the $\nu = 2$ quantum Hall resistance plateau we measure the breakdown current $I_c$, defined above, as a function of the magnetic field. Unlike the conventional QHE materials,11 the breakdown current in Fig. 3(a) continuously increases from zero to almost $500 \mu\text{A}$, far beyond $B_{\nu=2} \approx 9.5 \text{ T}$ calculated from the zero-field carrier density. This is a direct consequence of the exchange of carriers
between graphene and the donors in the “dead” layer, which keeps the $N = 0$ LL completely filled well past $B_{F=2}$.

The magnetic field range where the Fermi energy in SiC/G lies half-way between the $N = 0$ and 1 LLs determines the activation energy $\hbar \sqrt{2Tv/\lambda_B} \sim 1000$ K for the dissipative transport. For such a high activation energy, the low-temperature dissipative transport is most likely to proceed through the variable range hopping (VRH) between surface donors in SiC, involving virtual occupancy of the LL states in graphene to which they are weakly coupled. Indeed, as shown in the inset of Fig. 3(b), the temperature dependence of the conductivity $\sigma_{xx}$ measured at $B = 13$ T obeys an exp($-\sqrt{T/T^*}$) dependence typical of the VRH mechanism. The $T^*$ values determined from the measurements at different magnetic fields are plotted in the main panel of Fig. 3(b). The breakdown current rising with field to very large values [Fig. 3(a)] corresponds to $T^*$, reaching extremely large values in excess of $10^4$ K—at least an order of magnitude larger than that observed in GaAs\textsuperscript{28} and more recently in exfoliated graphene\textsuperscript{13,29}

In conclusion, we have studied the robust Hall resistance quantization in a large epitaxial graphene sample grown on SiC. We have observed the pinning of the $v = 2$ state, which is consistent with our picture of magnetic-field-dependent charge transfer between the SiC surface and graphene layer. Together with the large breakdown current this makes graphene on SiC the ideal system for high-precision resistance metrology.

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